

Molecular impurities interacting with a many-body environment: dynamics in Helium nanodroplets

G. Bighin and M. Lemeshko

Institute of Science and Technology Austria

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Quantum impurities

Definition: one (or a few particles) interacting with a many-body environment.

How are the properties of the particle modified by the interaction?

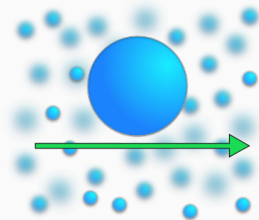
$\mathcal{O}(10^{23})$ degrees of freedom.



From impurities to quasiparticles

Structureless impurity: translational degrees of freedom/linear momentum exchange with the bath.

Most common cases: electron in a solid, atomic impurities in a BEC.



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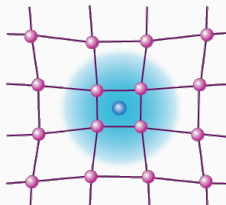


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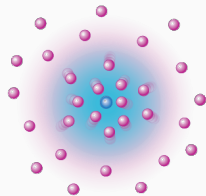


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From impurities to quasiparticles

Structureless impurity: translational

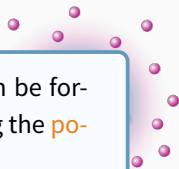
degrees of freedom
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atomic impurities in a BEC.

This scenario (with a bosonic bath) can be formalized in terms of **quasiparticles** using the **polaron** and the **Fröhlich** Hamiltonian.

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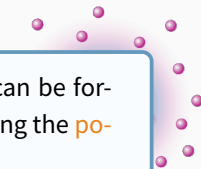
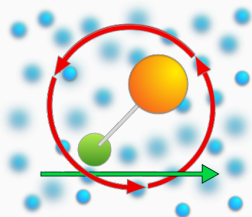


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Composite impurity: translational *and* internal (i.e. rotational) degrees of freedom/linear and angular momentum exchange.

From impurities to quasiparticles

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What about a **rotating particle**? Can there be a **rotating counterpart** of the **polaron quasiparticle**? The main difficulty: the **non-Abelian $SO(3)$ algebra** describing rotations.

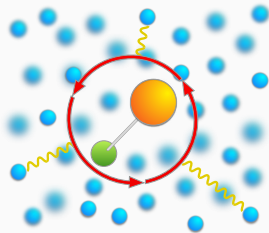
and internal
near and

The angulon

A composite impurity in a bosonic environment can be described by the angulon Hamiltonian^{1,2,3,4} (angular momentum basis: $\mathbf{k} \rightarrow \{k, \lambda, \mu\}$):

$$\hat{H} = \underbrace{B\hat{J}^2}_{\text{molecule}} + \underbrace{\sum_{k\lambda\mu} \omega_k \hat{b}_{k\lambda\mu}^\dagger \hat{b}_{k\lambda\mu}}_{\text{phonons}} + \underbrace{\sum_{k\lambda\mu} U_\lambda(k) \left[Y_{\lambda\mu}^*(\hat{\theta}, \hat{\phi}) \hat{b}_{k\lambda\mu}^\dagger + Y_{\lambda\mu}(\hat{\theta}, \hat{\phi}) \hat{b}_{k\lambda\mu} \right]}_{\text{molecule-phonon interaction}}$$

- Linear molecule.
- Derived rigorously for a molecule in a weakly-interacting BEC¹.
- Phenomenological model for a molecule in any kind of bosonic bath³.



¹R. Schmidt and M. Lemeshko, Phys. Rev. Lett. **114**, 203001 (2015).

²R. Schmidt and M. Lemeshko, Phys. Rev. X **6**, 011012 (2016).

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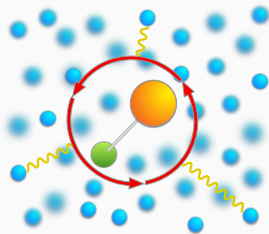
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$\lambda = 0$: spherically symmetric part.

$\lambda \geq 1$ anisotropic part.

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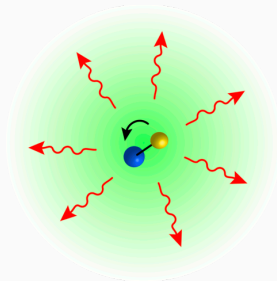
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Composite impurities: where to find them

Strong motivation for the study of composite impurities comes from many different fields. Composite impurities are realized as:

- **Ultracold molecules** and ions.

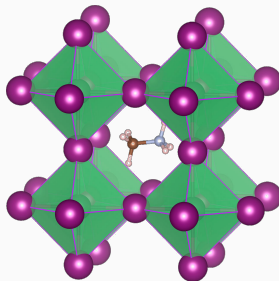


B. Midya, M. Tomza, R. Schmidt, and M. Lemeshko, Phys. Rev. A **94**, 041601(R) (2016).

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T. Chen et al., PNAS **114**, 7519 (2017).

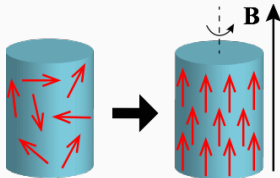
J. Lahnsteiner et al., Phys. Rev. B **94**, 214114 (2016).

Image from: C. Eames et al, Nat. Comm. **6**, 7497 (2015).

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- Angular momentum transfer from the electrons to a crystal lattice.



J.H. Mentink, M.I. Katsnelson, M. Leshenko, “Quantum many-body dynamics of the Einstein-de Haas effect”, arXiv:1802.01638

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- **Ultracold molecules** and ions.
- Rotating molecules inside a 'cage' in **perovskites**.
- Angular momentum transfer from the **electrons** to a **crystal lattice**.
- **Molecules** embedded into **helium nanodroplets**.

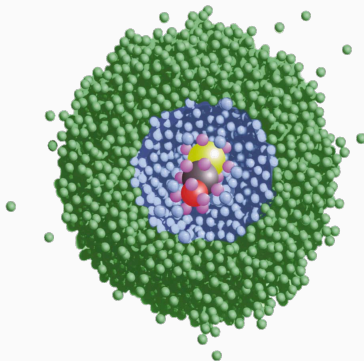
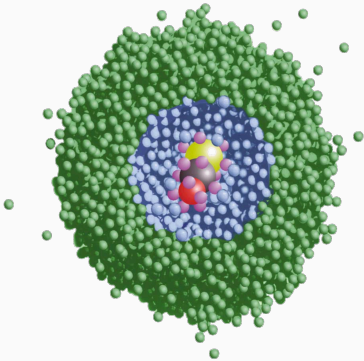


Image from: J. P. Toennies and A. F. Vilesov, *Angew. Chem. Int. Ed.* **43**, 2622 (2004).

Out-of-equilibrium dynamics of molecules in He nanodroplets

Dynamical alignment of molecules in He nanodroplets

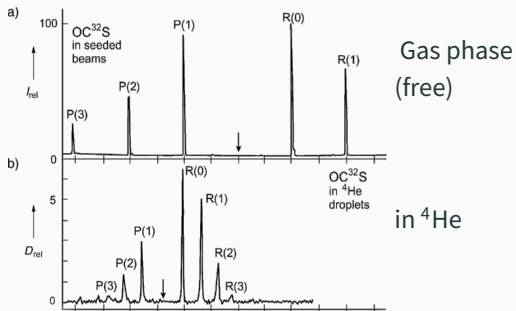
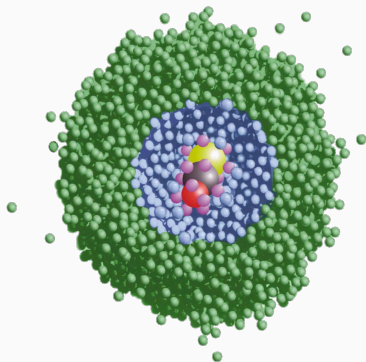
Molecules embedded into helium nanodroplets:



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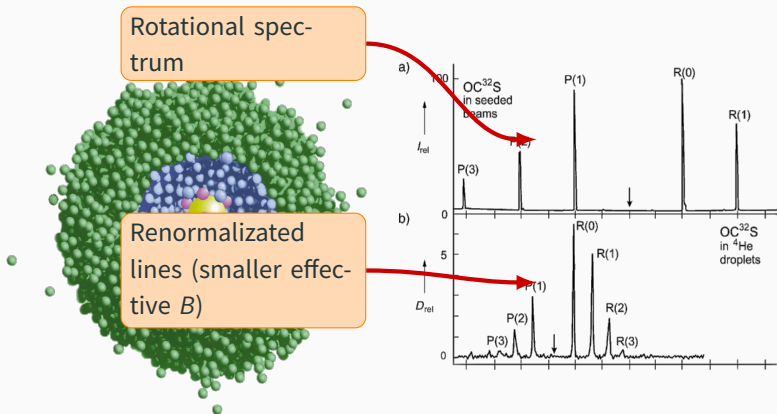
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Dynamical alignment of molecules in He nanodroplets

Dynamical alignment experiments:

- **Kick** pulse, aligning the molecule.
- **Probe** pulse, destroying the molecule.
- Fragments are imaged, reconstructing alignment as a function of time.
- Averaging over multiple realizations, and varying the time between the two pulses, one gets

$$\langle \cos^2 \hat{\theta}_{2D} \rangle (t)$$

with:

$$\cos^2 \hat{\theta}_{2D} \equiv \frac{\cos^2 \hat{\theta}}{\cos^2 \hat{\theta} + \sin^2 \hat{\theta} \sin^2 \hat{\phi}}$$

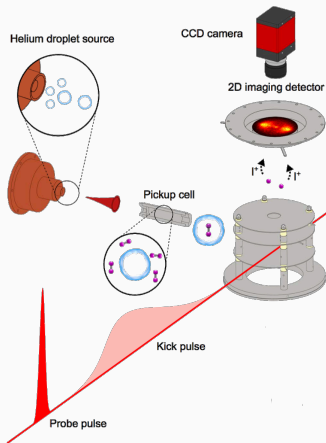


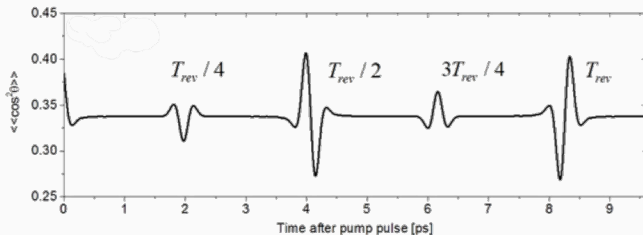
Image from B. Shepperson et al., Phys. Rev. Lett. **118**, 203203 (2017).

Dynamical alignment of molecules in He nanodroplets

Interaction of a **free molecule** with an off-resonant laser pulse

$$\hat{H} = B\hat{J}^2 - \frac{1}{4}\Delta\alpha E^2(t) \cos^2 \hat{\theta}$$

When acting on a **free molecule**, the laser excites in a short time many rotational states ($L \leftrightarrow L + 2$), creating a **rotational wave packet**:

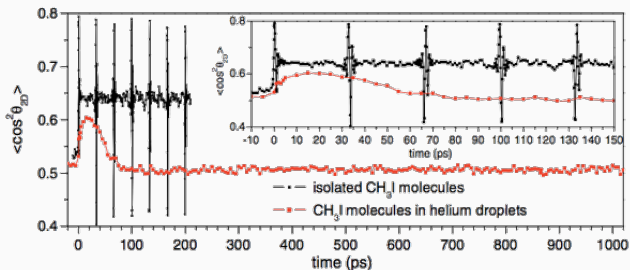


G. Kaya, Appl. Phys. B 6, 122 (2016).

Movie

Dynamical alignment of molecules in He nanodroplets

Effect of the environment is substantial: free molecule vs. **same molecule in He**.



Stapelfeldt group, Phys. Rev. Lett. **110**, 093002 (2013).

Not even a qualitative understanding. Monte Carlo?

- Strong coupling
- Out-of-equilibrium dynamics
- Finite temperature ($B \sim k_B T$)

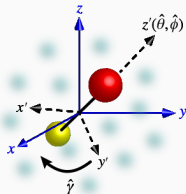
Canonical transformation

Bosons: laboratory frame (x, y, z)

Molecules: rotating frame (x', y', z')
defined by the Euler angles $(\hat{\phi}, \hat{\theta}, \hat{\gamma})$.

$$\hat{S} = e^{-i\hat{\phi} \otimes \hat{\Lambda}_z} e^{-i\hat{\theta} \otimes \hat{\Lambda}_y} e^{-i\hat{\gamma} \otimes \hat{\Lambda}_z}$$

where $\vec{\hat{\Lambda}} = \sum_{\mu\nu} b_{k\lambda\mu}^\dagger \vec{\sigma}_{\mu\nu} b_{k\lambda\nu}$ is the angular momentum of the bosons.



The \hat{S} transformation takes us to the molecular frame.

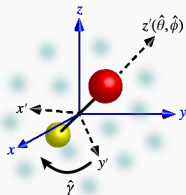
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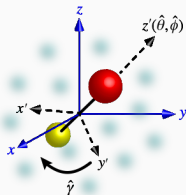
- **Macroscopic deformation** of the bath, exciting an infinite number of bosons (cf. Lee-Low-Pines for the polaron).
- Simplifies angular momentum algebra.
- Hamiltonian diagonalizable through a coherent state transformation in the $B \rightarrow 0$ limit. An expansion in bath excitations is a **strong coupling** expansion.

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The \hat{S} transformation takes us to the rotating frame.

- **Macroscopic** – Out-of-equilibrium dynamics
- **Simplified** – Finite temperature ($B \sim k_B T$)
- **Hamiltonian** – obtainable through a coherent state transformation in the $B \rightarrow 0$ limit. An expansion in bath excitations is a **strong coupling** expansion.

Dynamics: time-dependent variational Ansatz

We use a **time-dependent variational Ansatz**:

$$|\psi\rangle = g_{LM}(t) |0\rangle_{\text{bos}} |LM0\rangle + \sum_{k\lambda n} \alpha_{k\lambda n}^{LM}(t) b_{k\lambda n}^\dagger |0\rangle_{\text{bos}} |LMn\rangle$$

Lagrangian on the variational manifold defined by $|\psi\rangle$:

$$\mathcal{L}_{T=0} = \langle \psi | i\partial_t - \hat{\mathcal{H}} | \psi \rangle$$

Euler-Lagrange equations of motion:

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{x}_i} - \frac{\partial \mathcal{L}}{\partial x_i} = 0$$

where $x_i = \{g_{LM}, \alpha_{k\lambda n}\}$.

$$\begin{cases} \dot{g}_{LM}(t) = \dots \\ \dot{\alpha}_{k\lambda n}^{LM}(t) = \dots \end{cases}$$

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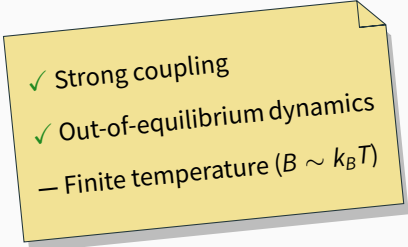
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- 
- ✓ Strong coupling
 - ✓ Out-of-equilibrium dynamics
 - Finite temperature ($B \sim k_B T$)

Finite-temperature dynamics

For the **impurity**: average over a statistical ensemble, with weights $W_L \propto \exp(-\beta E_L)$.

For the **bath**: defining the ‘Chevy operator’

$$\hat{O} = g_{LM}(t) |LM0\rangle \mathbb{1} + \sum_{k\lambda n} \alpha_{k\lambda n}^{LM}(t) |LMn\rangle \hat{b}_{k\lambda n}^\dagger$$

at $T = 0$ the Lagrangian is

$$\mathcal{L}_{T=0} = \langle 0 | \hat{O}^\dagger (i\partial_t - \hat{\mathcal{H}}) \hat{O} | 0 \rangle_{\text{bos}} ,$$

suggesting that at **finite temperature**

$$\mathcal{L}_T = \text{Tr} \left[\rho_0 \hat{O}^\dagger (i\partial_t - \hat{\mathcal{H}}) \hat{O} \right]$$

where ρ_0 is the **density matrix** for the medium.

[1] A. R. DeAngelis and G. Gatoff, Phys. Rev. C **43**, 2747 (1991).

[2] W.E. Liu, J. Levinsen, M. M. Parish, “Variational approach for impurity dynamics at finite temperature”, arXiv:1805.10013

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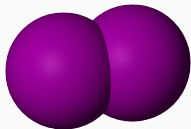
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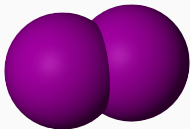
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Theory vs. experiments: I_2



Comparison of the theory with preliminary experimental data from Stapelfeldt group, Aarhus University, for different molecules: I_2 .

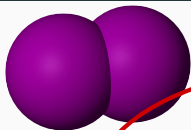
Theory vs. experiments: I_2



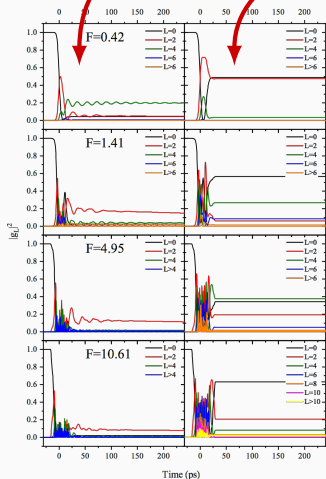
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Which rotational states are populated as the laser is switched on, and after?

Theory vs. experiments: l_2



Comparison of the theory with preliminary experiment in Helium University, Free molecule: l_2 .

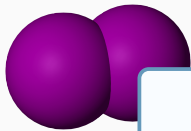


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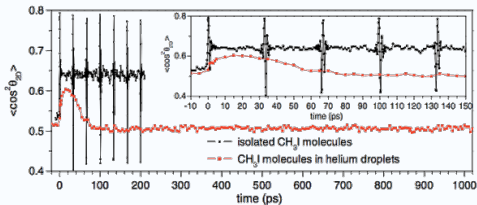
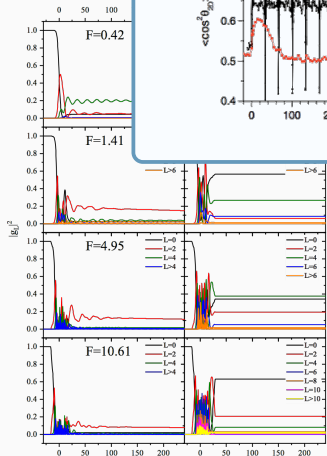
Free case: the angular momentum goes to the molecule.

In a Helium droplet: the angular momentum goes to the molecule *and* to the bath.

Theory vs. experiments: l_2



Comparison of the theory with preliminary experimental data from Stapelfeldt group, Aarhus

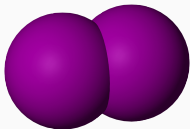


re
switched

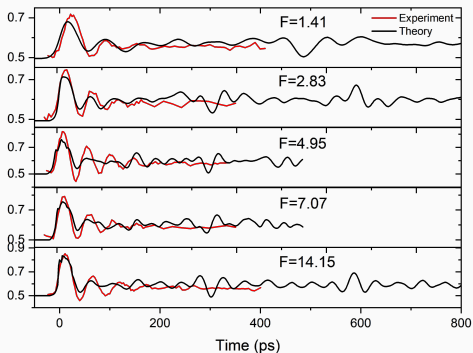
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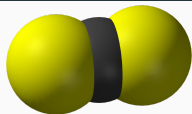
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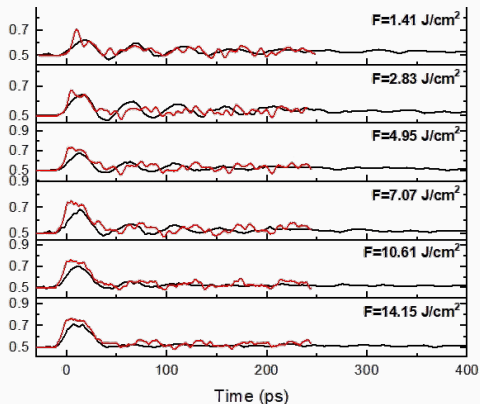
$$\langle \cos^2 \hat{\theta}_{2D} \rangle (t)$$

Laser fluence F
measured in J/cm^2

Theory vs. experiments: CS_2



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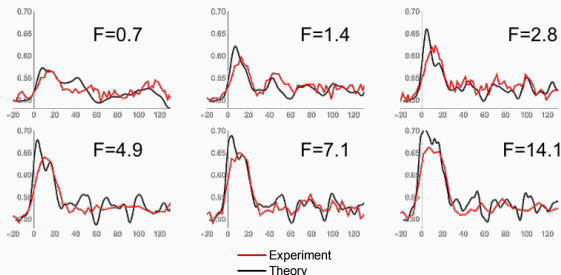


$$\langle \cos^2 \hat{\theta}_{2D} \rangle (t)$$

Theory vs. experiments: OCS



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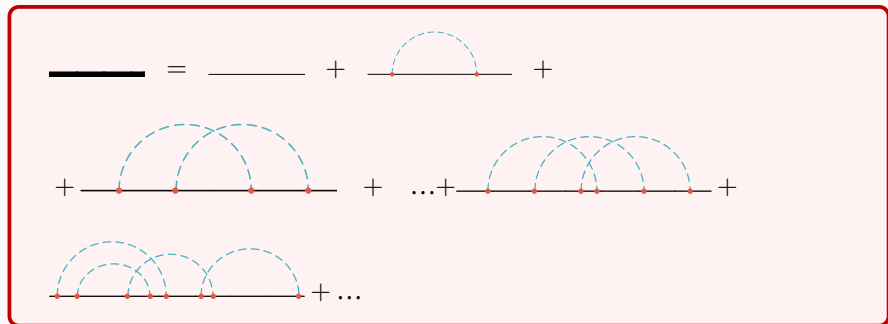


$$\langle \cos^2 \hat{\theta}_{2D} \rangle (t)$$

Laser fluence F
measured in J/cm^2 ,
time measured in ps .

Diagrammatic Monte Carlo

More numerical approach: **DiagMC**, sampling all diagrams in a stochastic way.



How do we describe angular momentum redistribution in terms of diagrams?
How does the configuration space look like?

Connecting DiagMC and the theory of molecular simulations!

Conclusions

- The **angulon quasiparticle**: a quantum rotor dressed by a field of many-body excitations.
- Canonical transformation and **finite-temperature** variational Ansatz.
- **Out-of-equilibrium dynamics** of molecules in He nanodroplets can be interpreted in terms of angulons.



Institute of Science and Technology

Lemeshko group @ IST Austria:



Misha
Lemeshko

Dynamics in He



Enderalp
Yakaboylu



Xiang Li



Igor
Cherepanov



Wojciech
Rządkowski

Collaborators:



Henrik
Stapelfeldt
(Aarhus)

Dynamical
alignment ex-
periments



Thank you for your attention.



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Backup slide # 1

Free rotor propagator

$$G_{0,\lambda}(E) = \frac{1}{E - B\lambda(\lambda + 1) + i\delta}$$

Interaction propagator

$$\chi_\lambda(E) = \sum_k \frac{|U_\lambda(k)|^2}{E - \omega_k + i\delta}$$

